

Class 3 – Photonic Crystals (October 8, 2010)

Assignment

In this assignment, we will develop a basic framework to describe the electromagnetic eigenfunctions in a one-dimensional photonic crystal. We will derive the formation of a photonic bandgap: a range of frequencies for which light can not propagate normal to the crystal planes.

- a) Starting from Maxwell's equations in the absence of free charges and currents, derive the wave equation

$$\frac{1}{\epsilon} \nabla \times (\nabla \times \mathbf{E}) = \left(\frac{\omega}{c} \right)^2 \mathbf{E}. \quad (1)$$

Equation (1) describes an eigenvalue problem for the electric field. The solution can be written as an expansion in plane waves through a Fourier transform:

$$\mathbf{E}(\mathbf{x}) = \int d^3k \mathbf{A}(\mathbf{k}) e^{-i\mathbf{k} \cdot \mathbf{x}}, \quad (2)$$

where $\mathbf{A}(\mathbf{k})$ is the amplitude of a plane wave with wave vector \mathbf{k} .

In a photonic crystal, the dielectric constant is periodic;

$$\epsilon(\mathbf{x}) = \epsilon(\mathbf{x} + \mathbf{a}), \quad (3)$$

where \mathbf{a} is an arbitrary lattice vector. For simplicity, in this assignment we will consider a one-dimensional photonic crystal (for example a stack of alternating layers of two different dielectrics), that is periodic in the z direction with period a . The reciprocal lattice vectors \mathbf{G} of the crystal are defined as

$$\mathbf{G} = n g \hat{\mathbf{z}} = n \frac{2\pi}{a} \hat{\mathbf{z}}, \quad n = 0, \pm 1, \pm 2, \dots \quad (4)$$

- b) Show that

$$\mathbf{k} \times (\mathbf{k} \times \mathbf{A}(\mathbf{k})) + \left(\frac{\omega}{c} \right)^2 \sum_{\mathbf{G}} \beta_{\mathbf{G}} \mathbf{A}(\mathbf{k} - \mathbf{G}) = 0, \quad (5)$$

where $\beta_{\mathbf{G}}$ are the Fourier coefficients of the dielectric function, $\epsilon(\mathbf{x})$.

We now want to understand Bloch's theorem. Bloch's theorem states that in a periodic medium with period a , eigenfunctions of the electric field can be described by

$$\mathbf{E}(\mathbf{x}) = \mathbf{E}_{\mathbf{K}}(\mathbf{x}) e^{-i\mathbf{K} \cdot \mathbf{x}}, \quad (6)$$

where $\mathbf{E}_{\mathbf{K}}$ is periodic; $\mathbf{E}_{\mathbf{K}}(\mathbf{x}) = \mathbf{E}_{\mathbf{K}}(\mathbf{x} + \mathbf{a})$. It depends on the wavevector \mathbf{K} , which is known as the Bloch wavevector.

- c) Consider eq. (5). The solutions to eq. (5) should be written in the form of a plane wave expansion. If the plane wave with a particular wavevector \mathbf{K} appears in the eigensolution (2), explain which other plane waves are part of that solution. Using this information and the plane wave expansion of the field (eq. (2)), write $\mathbf{E}_{\mathbf{K}}(\mathbf{x})$ in terms of the coefficients \mathbf{A} . Explain that you have now derived Bloch's theorem.

In the rest of this assignment we will assume wave propagation in the z direction. The photonic crystal is homogeneous in the x and y directions, and $\varepsilon(z)$ is always real and positive. To describe wave propagation in the crystal, it is necessary to find a dispersion relation that relates ω and K .

- d) Solving equation (5) can yield both the dispersion relation and the eigenfunctions $\mathbf{E}_{\mathbf{K}}$. Equation (5) is in principle an infinite set of coupled equations, but in practice only a few equations are needed to provide an approximate result. Expand the dielectric function only up to the first three terms; β_0 , β_{-1} , and β_1 (where the subscript denotes n in eq. (4)). Note that $\mathbf{k} \cdot \mathbf{A} = 0$. Near the Bragg condition, the dominant terms in the equations following from eq. (5) are only those with $A(K)$ and $A(K - g)$. Show that the dispersion relation near the Bragg condition is approximated by:

$$(\omega/c)^4 (\beta_0^2 - \beta_1 \beta_{-1}) - (\omega/c)^2 \beta_0 (K^2 + (K - g)^2) + K^2 (K - g)^2 = 0 \quad (7)$$

Show that when the Bragg condition is satisfied exactly, two roots for ω^2 may be obtained:

$$\omega_{\pm}^2 = \frac{\pi^2 c^2}{a^2 (\beta_0 \pm |\beta_1|)} \quad (8)$$

The frequencies between ω_+ and ω_- are within the bandgap that has arisen from the coupling between the wave components $A(K)$ and $A(K - g)$.

- e) Give an expression for K for the center frequency of the stop gap (at $\omega = \pi c / (a \sqrt{\beta_0})$). Use an approximation assuming K is very close (but not equal) to $g/2$. Describe qualitatively what happens to a wave that is incident *from outside* onto a finite-sized photonic crystal at this frequency.
- f) The fractional bandgap (the frequency width of the stop gap divided by the center frequency in the stop gap) is a good measure of the 'photonic strength' of the photonic crystal, i.e., the coupling strength between $A(K)$ and $A(K - g)$ at the Bragg condition. Consider a one-dimensional photonic crystal that consists of alternating layers of equal thickness of two materials with refractive indices of 3.0 and 3.4. What is the magnitude of the fractional bandgap?

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